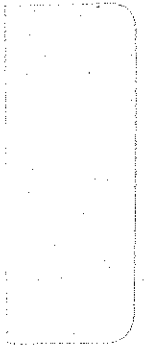


# REPRINT



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SOME NONLINEAR FILTERING PROBLEMS ARISING IN RECURSIVE IDENTIFICATION

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Abstract: In this paper, we outline an intrinsic formulation of the identification problem of linear system theory. The nonlinear filtering problems which appear in this way essentially fall into four distinct classes, distinguished by their estimation algebra. In principle, it is possible to explicitly solve the identification problem in the 'hyperbolic cases' using classical methods from the theory of partial differential equations. This is illustrated by an example which indicates the required sufficient statistics for solving the identification problem.

1. INTRODUCTION

Consider the linear stochastic differential system,

$$dx_t = A(p)x_t dt + b_1 u_t dt + b_2 dw_t \tag{1.1}$$

$$dy_t = \langle q, x_t \rangle dt + dv_t$$

where  $u_t$  denotes a known input function,  $\{w_t\}$  and  $\{v_t\}$  are independent Brownian motion processes,  $\{x_t\}$  and  $\{y_t\}$  are respectively the state and measured output processes. For reasons of identifiability we let,

$$A(p) = \begin{bmatrix} 0 & 1 & & & \\ & & 1 & & \\ & & & \ddots & \\ & & & & 1 \\ -p_1 & -p_2 & \cdots & -p_n & \end{bmatrix}$$

the rational canonical form associated with  $p=(p_1, \dots, p_n)' \in \mathbb{R}^n$  and we let,  $q=(q_1, \dots, q_n)' \in \mathbb{R}^n$ . The vectors  $b_1$  and  $b_2$  are known and fixed. When  $p$  and  $q$  are known, the state-estimation problem for (1.1) has the well-known solution -- the Kalman-Bucy filter. By the identification problem we shall mean the problem of jointly estimating the state and the parameters -- in other words it is the nonlinear filtering problem for the extended system with state  $z_t=(x_t, p_t, q_t)$  defined by,

$$\begin{aligned} dx_t &= A(p_t)x_t dt + b_1 u_t dt + b_2 dw_t \\ dp_t &= 0 \end{aligned} \tag{1.3}$$

$$\begin{aligned} dq_t &= 0, \\ dy_t &= \langle q_t, x_t \rangle dt + dv_t. \end{aligned} \tag{1.4}$$

More precisely, in solving the identification problem one seeks a solution to the Kushner-Stratonovitch eqn, [1] satisfied by the conditional density  $p(t, z)$  given the observations  $y_s, 0 \leq s \leq t$ . Although this point of view goes back to Kushner [2], progress along these lines has been impeded by the nonlinearity of the Kushner equation.

More recently, it has been recognized [3,4,5] that an understanding of the evolution equation satisfied by the so-called unnormalized conditional density,  $\psi(t, z)$  is essential for further progress in nonlinear filtering theory. Knowing  $\psi(t, z)$ , the conditional density is determined by the normalization

$$p(t, x) = \psi(t, z) / \int \psi(t, z) dz. \tag{1.5}$$

In the general situation when  $\{z_t\}$  is a diffusion process with observation (semimartingale)  $\{y_t\}$  of the form

$$dy_t = h(z_t) dt + dv_t \tag{1.6}$$

it is known that,  $\psi(t, z)$  satisfies a linear stochastic partial differential equation (Mortenson-Duncan-Zakai equation) of the Ito type,

$$d\psi(t, z) = \mathcal{L}\psi(t, z) dt + h(z) \psi(t, z) dy_t \tag{1.7}$$

where  $\mathcal{L}$  is the Kolmogorov forward operator associated with the diffusion  $\{z_t\}$ . (See the expository paper by Davis and Marcus in these proceedings). Now, as regards questions related to the complexity of a nonlinear filtering problem geometric ideas play a crucial role and one looks at the Stratonovitch version of (1.7) written formally as,

$$\partial\psi/\partial t(t,z)=[\mathcal{L}^{-1/2}h^2(z)]\psi(t,z)+h(z)dy/dt. \quad (1.8)$$

In particular, the operator Lie algebra  $G$  generated by  $\mathcal{L}^{-1/2}h^2$  and  $h$ , known as the estimation algebra [3], has been emphasized by Brockett, Mitter, Ocone and others as an object of central interest. In what follows we use the estimation algebra to classify identification problems and investigate special cases.

2. ESTIMATION ALGEBRAS:

The fact that the estimation algebra is invariant under change of coordinates [6], makes it useful as a classifying tool. In particular the choice of canonical forms is not crucial. Essentially there are four cases:

Case (1)  $b_2=0; b_1=0$ . Then  $\mathcal{L}=-\langle Ax, \partial/\partial x \rangle - \text{tr}(A)$ . Define  $A_0 = \mathcal{L}^{-1/2}\langle q, x \rangle^2; A_1 = \langle q, x \rangle$ . Then the estimation algebra  $G = \{A_0, A_1\}$  L.A. is defined by the structure equations,  $[A_0, A_j] = A_{j+1}; [A_j, A_k] = 0, k, j \geq 1$  where,  $A_j = (-1)^{j-1}\langle q, A^{j-1}x \rangle, j=1, 2, \dots$ . We have a sequence of abelian ideals  $G_n = \text{span}\{A_j; j=n, n+1, \dots\}, n \geq 1$ , with finite codimension (a feature of potential value in connection with approximation problems).

Case (2)  $b_1 = e_n$  say;  $b_2=0$ . The presence of deterministic inputs does not alter the structure of the estimation algebra. Note  $\mathcal{L} = -\langle Ax, \partial/\partial x \rangle - \text{tr}(A) - u_t \langle b_1, \partial/\partial x \rangle$ . Define  $A_0 = \mathcal{L}^{-1/2}\langle q, x \rangle^2$  and  $A_j = (-1)^{j-1}\langle q, A^{j-1}x \rangle + (-1)^{j-1}\langle q, A^{j-2}b_1 \rangle u_t, j \geq 2$ . Then once again  $[A_0, A_j] = A_{j+1}; [A_j, A_k] = 0; \text{ for } j, k \geq 1$ .

Case (3)  $b_1 = e_n; b_2=0$  and the parameters  $p$  are known.  $A_0, A_1, \dots$  etc are as in the previous case. But by the Cayley-Hamilton theorem  $A^n(p) = \sum_{i=1}^n p_i A^{i-1}(p)$ . Hence the operators  $A_k$  for  $k > n+1$  are

linearly dependent on the operators  $A_j$  for  $j \leq n+1$ . Then the estimation algebra  $G = \{A_0, A_1\}$  L.A. is finite dimensional. In fact using tensor products it can be shown that the underlying filtering problem is linear.

Case (4)  $b_2 \neq 0$ . The presence of driving noise drastically alters the structure of the Lie algebra. The general situation is not unlike the example below:

$$\begin{aligned} dx_t &= \alpha x_t dt + dw_t \\ da_t &= 0 \\ dy_t &= x_t dt + dv_t \end{aligned}$$

$$\mathcal{L} = -\alpha x \partial/\partial x - \alpha + \frac{1}{2} \partial^2/\partial x^2; A_0 = \mathcal{L}^{-1/2} x^2; A_1 = x.$$

Define  $A_{2n+1} = (\alpha^2 + 1)^n x; A_{2n+2} = (\alpha^2 + 1)^n (\partial/\partial x - \alpha x)$  with  $n=0, 1, 2, \dots$ . Also let  $B_k = -(\alpha^2 + 1)^k, k=0, 1, 2, \dots$ . Then the structure equations are,

$$\begin{aligned} [A_0, A_{2n+1}] &= A_{2n+2}; [A_0, A_{2n+2}] = A_{2n+3} [A_{2n+1}, A_{2m+1}] = 0; \\ [A_{2n+2}, A_{2m+2}] &= 0; [A_{2n+1}, A_{2m+2}] = B_{m+n}. [B_j, A_k] = 0; [B_j, B_k] = 0. \end{aligned}$$

It is possible to write down filtrations of  $\bar{G}$  by sequences of ideals as before. In fact in each case above the algebra  $\bar{G}$  is a profinite dimensional filtered Lie algebra (see Hazewinkel-Marcus [10]). All the known nonlinear filtering problems that admit finite dimensionally computable statistics have Lie algebras of this type.

It is however important to note that the identification problem in our formulation is tractable in the cases (1), (2), (3) above where there is no driving noise. In these cases, the Stratonovitch form of the evolution equation for the unnormalized conditional density is given by,

$$\partial\psi/\partial t = (A_0 + A_1\dot{y})\psi \quad (2.1)$$

where,  $A_0 = -\langle Ax, \partial/\partial x \rangle - \text{tr}(A) - u_t \langle b_1, \partial/\partial x \rangle - \frac{1}{2} \langle q, x \rangle^2$ , and  $A_1 = \langle q, x \rangle$ . In principle, equation (2.1) can be solved by the method of characteristics. See below.

### 3. AN EXAMPLE

Consider the special case of (1.3)-(1.4) given by

$$\begin{aligned} dx_t &= -\alpha x_t dt + u_t dt \\ dy_t &= x_t dt + dv_t \end{aligned} \quad (3.1)$$

Then (2.1) reduces to,

$$\partial\psi/\partial t = (\alpha x - u_t) \partial\psi/\partial x + (\alpha - x^2/2 + x\dot{y})\psi. \quad (3.2)$$

Let the initial condition be given by  $\psi(t, x, \alpha) |_{t=0} = \psi_0(x, \alpha)$ . In the 4-dimensional  $(t, x, \alpha, z)$  space, we want to pass an integral hyper-surface  $S: z = \psi(t, x, \alpha)$  of the equation (3.2) through the 2-dimensional manifold  $\Gamma$  (Cauchy data) given parametrically by  $x = s_1, \alpha = s_2, t = 0, z = \psi_0(s_1, s_2)$ . The characteristics passing through points  $(s_1, s_2)$  in  $\Gamma$  sweep out  $S$  and are given by the system of (characteristic) differential equations:

$$\begin{aligned} dx/d\tau &= -(\alpha x - u_t) \\ d\alpha/d\tau &= 0 \\ dt/d\tau &= 1 \\ dz/d\tau &= (\alpha - x^2/2 + x\dot{y}_t)z \end{aligned} \quad (3.3)$$

Solving (3.3), we obtain a parametric representation of the characteristic curves;  $\alpha = s_2; t = \tau; x = X(s_1, s_2, \tau) = e^{-s_2\tau} s_1 + \int_0^\tau e^{-s_2(\tau-\sigma)} u_\sigma d\sigma$  and finally,

$$z = \psi_0(s_1, s_2) \cdot \exp(s_2\tau) \cdot \exp\left(\int_0^\tau X(s_1, s_2, \sigma) \dot{y}_\sigma d\sigma - \frac{1}{2} \int_0^\tau X^2(s_1, s_2, \sigma) d\sigma\right). \quad (3.4)$$

Equation (3.4) for  $z$  is nothing but a parametric representa-

tion of the solution  $\psi$  we are seeking. It is easy to see that given  $x, t$ , and  $\alpha$  the parameters  $s_1, s_2$  and  $\tau$  can be eliminated and

$$X(s_1, s_2, \sigma) = e^{\alpha(t-\sigma)} x - \int_{\sigma}^t e^{-\alpha(\sigma-\theta)} u_{\theta} \cdot d\theta. \quad (3.5)$$

Substitution in (3.4) gives the explicit representation of  $\psi(t, x, \alpha)$  for given input and output functions. The last exponential factor in (3.4) corresponds to the well-known likelihood ratio formula ([7], [8]).

#### 4. SUFFICIENT STATISTICS

In Eqn. (3.4), only the term  $\int_0^t X(s_1, s_2, \sigma) \dot{y}_{\sigma} d\sigma$  inside the exponential depends on measured outputs and explicitly,

$$\begin{aligned} \int_0^t X(s_1, s_2, \sigma) \dot{y}_{\sigma} d\sigma \\ = x \sum_{k=0}^{\infty} \alpha^k \int_0^t \frac{(t-\sigma)^k}{k!} dy_{\sigma} - \sum_{k=0}^{\infty} (-\alpha)^k \int_0^t \gamma_k(t, \sigma) dy_{\sigma} \end{aligned} \quad (4.1)$$

where  $\gamma_k(t, \sigma) = \int_{\sigma}^t (\sigma-\theta)^k / k! u_{\theta} d\theta$ . The two sequences

$$(a) \quad \beta_k(t) = \int_0^t (t-\sigma)^k / k! dy_{\sigma} \quad k = 0, 1, 2 \dots$$

and

$$(b) \quad \omega_k(t) = \int_0^t \gamma_k(t, \sigma) dy_{\sigma} \quad k = 0, 1, 2, \dots$$

may be viewed as sufficient statistics for the problem. Each  $\beta_k$  can be computed as the output of a finite dimensional system driven by  $y_t$ . The same holds true for the  $\omega_k$ 's. We must mention that the statistics  $\omega_k$  are similar to the sufficient statistics determining the likelihood ratio given by Giorgio Picci [9].

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